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Physica B 241–243 (1998) 702–703

PHYSICA B

Neutron diffraction studies of Pr ordering in $\text{PrBa}_2\text{Cu}_4\text{O}_8$

Y.-C. Lin^a, S.Y. Wu^a, W.-H. Li^{a,*}, K.C. Lee^a, J.W. Lynn^b, C.W. Lin^c, J.-Y. Lin^c, H.D. Yang^c^a Department of Physics, National Central University, Chung-Li, 32054, Taiwan^b NIST Center for Neutron Research, NIST, Gaithersburg, MD 20899, USA^c Department of Physics, National Sun Yat-Sen University, Kaohsiung, 80424, Taiwan

Abstract

Neutron diffraction and AC susceptibility measurements have been performed to study the magnetic ordering of the Pr spins in $\text{PrBa}_2\text{Cu}_4\text{O}_8$. An anomaly at 15 K, which we believe is associated with the ordering of the Pr spins, is clearly seen in the temperature dependence of the AC susceptibility, $\chi'(T)$. The coupling between the Pr spins, as revealed by neutrons, is antiferromagnetic in nature, with a spin structure may be characterized by the $\{\frac{1}{2} \frac{1}{2} 1\}$ wave vector. Three-dimensional coupling between the Pr spins is observed, despite the appearance of exact cancellations of the coupling between the adjacent layers along the c -axis due to geometric symmetry. © 1998 Elsevier Science B.V. All rights reserved.

Keywords: High- T_c compound; Neutron diffraction; Magnetic structure

The Pr-based 123-cuprate has distinguished itself from other high- T_c cuprates by its special properties [1,2]. In continuously exploring the mysteries of Pr cuprates, we report neutron diffraction and AC susceptibility studies of the magnetic ordering of the Pr spins in a $\text{PrBa}_2\text{Cu}_4\text{O}_8$ (Pr124) compound. The compound is synthesized at ambient oxygen pressure by the nitric pyrolysis method [3,4].

Fig. 1 shows the temperature dependence of the in-phase component of the AC susceptibility. The solid line shown is a fit of the data obtained between 25 and 80 K to the Curie–Weiss law. Clearly, below 15 K, the data depart from the fitted line. The origin of this anomaly is believed to be associated with the Pr spin ordering.

*Corresponding author. Fax: + 886 3 4251175; e-mail: whli@joule.phy.ncu.edu.tw.

Neutron diffraction measurements were performed on the US NIST BT-2 triple-axis spectrometer, operated in double-axis mode with pyrolytic graphite monochromator and filter, with an incident wavelength of 2.351 Å and angular collimations 40'–48'–48'. The magnetic peaks observed at 6 K are shown in Fig. 2, where the data taken at 25 K have been subtracted to isolate the magnetic signal. Two magnetic peaks were observed, with widths broader than the instrumental resolution. Due to a long lattice constant for the c -axis, the expected separation in 2θ between the $\{\frac{1}{2} \frac{1}{2} 0\}$ and $\{\frac{1}{2} \frac{1}{2} 1\}$ peaks is only 0.8°, which is beyond our instrumental resolution. If we fit the data assuming the peak at 25° contains two peaks, the fitted peak positions are at the expected positions for the $\{\frac{1}{2} \frac{1}{2} 0\}$, $\{\frac{1}{2} \frac{1}{2} 1\}$ and $\{\frac{1}{2} \frac{1}{2} 2\}$ reflections and all three peaks are resolution limited. We hence believe that

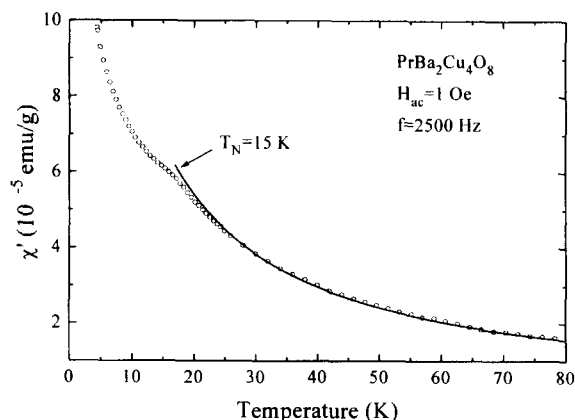


Fig. 1. Temperature dependence of the AC susceptibility revealing an anomaly at 15 K.

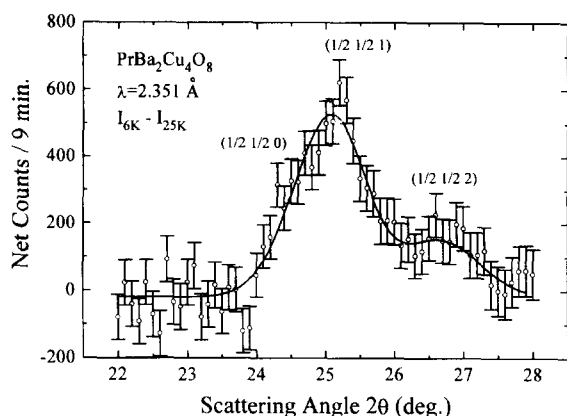


Fig. 2. The magnetic diffraction peaks observed at 6 K.

three magnetic peaks, which may be indexed as the $\{\frac{1}{2} \frac{1}{2} 0\}$, $\{\frac{1}{2} \frac{1}{2} 1\}$, and $\{\frac{1}{2} \frac{1}{2} 2\}$ reflections, were observed. The solid curve shown in Fig. 2 is a fit of the data to three Gaussian resolution functions. By comparing the $\{\frac{1}{2} \frac{1}{2} 1\}$ magnetic intensity observed at 6 K with that of $\{006\}$, we obtain a saturated

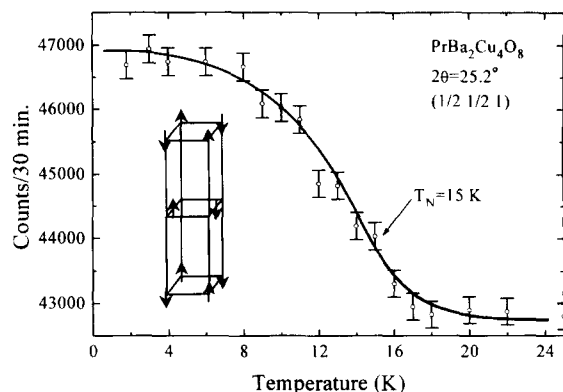


Fig. 3. Variation of the $\{\frac{1}{2} \frac{1}{2} 1\}$ peak intensity. The inset shows the proposed spin structure.

moment of $\langle \mu_z \rangle = 0.61(4)\mu_B$ where the moment points in the c -axis direction.

In Fig. 3, we show the variation of the $\{\frac{1}{2} \frac{1}{2} 1\}$ peak intensity with temperature. The ordering temperature, as determined by the inflection point, is $T_N = 15$ K, which matches the temperature at which $\chi'(T)$ shows an anomaly. Shown as an inset in Fig. 3 is the proposed Pr spin structure, in which the adjacent spins in the a - b plane are antiparallel and the coupling between the adjacent layers along the c -axis are canceled due to geometric symmetry.

This work was supported by the National Science Council of the Republic of China under Grant No. NSC 86-2112-M-008-029 and NSC 86-2112-M-110-013.

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